

- by this scenario, we have a repeated bifurcation sequence
→ turbulence.

- note: $l > 0$ *Bernard convection*
 $l > 0$ *Taylor cells*.
 $l < 0$ *Unbounded jets + shear layers*.
 $l < 0$ *Boundary layers, channel flows, pipe flows*.

- Landau's Theory comes from implied assumption that
that there is λ_1 unique and discrete most unstable
mode.
 - applicable to a bounded region where we expect linear eigenvalue problem with normal modes which are both discrete and complete,
 - not applicable in cases where normal modes depend continuously on wave number.

Bifurcations

Example: $\frac{1}{R} u'' + \sin u = 0$; $0 \leq y \leq \pi$
 $R > 0$

- one solution is $u=0$.

- linear theory for small displacements from $u=0$.

$$\frac{1}{R} u'' + u = 0$$

sol: $u = A \sin \sqrt{\frac{1}{R}} y + B \cos \sqrt{\frac{1}{R}} y$

apply BC : @ $y=0, u=0 \rightarrow B=0$

$$@ y=\pi, u=0 \rightarrow \sqrt{R} = n = \text{integer}_{1, 2, \dots}$$

or $R = n^2$



small amplitude expansion :

$$u = \epsilon u_1 + \epsilon^2 u_2 + \epsilon^3 u_3 + \dots$$

$$R = n^2 + \epsilon R_1 + \epsilon^2 R_2 + \dots$$

ϵ is a small parameter which is a measure of amplitude

$$\text{then } u'' + R \sin u = 0$$

$$\text{gives } u'' + R \left\{ u - \frac{u^3}{3!} + \dots \right\} = 0$$

or,

$$\epsilon u_1'' + \epsilon^2 u_2'' + \epsilon^3 u_3'' + \dots + \left\{ R^2 + \epsilon R_1 + \epsilon^2 R_2 + \dots \right\}$$

$$\cdot \left\{ \epsilon u_1 + \epsilon^2 u_2 + \epsilon^3 u_3 - \frac{\epsilon^3 u_1^3}{3!} \dots \right\}$$

equate like powers of ϵ ,

$$\boxed{\left\{ \begin{array}{l} u_1'' + u_1 = 0 \\ u_2'' + u_2 = -R_1 u_1 \\ u_3'' + u_3 = R_1 u_2 - R_2 u_1 + \frac{u_1^3}{3!} \end{array} \right.} \quad \begin{array}{l} \text{BC: } u_i(0) = 0 \\ u_i(\pi) = 0 \end{array} \quad i=1,2,3$$

solutions:

$$\textcircled{1} \quad u_1 = A \sin y$$

$$\textcircled{2} \quad \text{then } u_2'' + u_2 = -R_1 A \sin y$$

$$u_2 = \alpha \sin y + \beta \cos y + \cancel{\gamma y \cos y}$$

$$\text{BC: } u_2(0) = u_2(\pi) = 0 \quad \Rightarrow \quad \beta = 0$$

$$\downarrow \\ \beta = 0$$

$$\rightarrow -2R_1 A$$

$$\Downarrow \\ 0 = -2\cancel{(\beta)} A \pi \cos \pi$$

$$\therefore R_1 = 0$$

\textcircled{3} next eq. becomes.

$$u_3'' + u_3 = -R_2 u_1 + \frac{u_1^3}{3!} = R_2 A \sin y$$

$$+ \frac{1}{3!} (A^3 \sin^3 y)$$

$$= R_2 A \sin y + \frac{A^3}{3!} \sin y + \frac{A^3}{3!} \sin 3y$$

$$\text{BC: } u_3(0) = u_3(\pi) = 0$$

$$\text{solvability cond: } -R_2 A + \frac{A^3}{3!} = 0$$

$$\text{if } A \neq 0 \text{ then } R_2 = \frac{A^2}{3!}$$

checking what we have:

$$u = \epsilon A \sin y + \epsilon^2 u_2 + \epsilon^3 u_3 + \dots$$

$$R-1 = \epsilon R_1 + \epsilon^2 R_2 + \dots = k \epsilon^2 A^2$$

$$|u| = \epsilon A + \dots \text{ H.o.T}$$

$$R-1 = k(\epsilon A)^2 + \dots = k |u|^2$$

