

Recall:

Lem: $\begin{matrix} V \\ \downarrow \\ X \end{matrix}$ vector ball. Then $\hat{\Omega}(x) := \Gamma(S^*(V) \otimes \Lambda^*(V^\vee[1]))$
 with differential $1 \otimes \xi[1] \xrightarrow{d} \xi \otimes 1 \xrightarrow{d} 0$
 is a free resolution of $\Omega(x)$ by free $\Omega(V) = \Gamma(S^*(V))$ -modules.

Note: $S^*(V) \otimes \Lambda^*(V^\vee[1])$ is a cochain complex of vector balls.

In particular, for each $x \in X$ it suffices to show

$$H^k(S^*(V_x) \otimes \Lambda^*(V_x^\vee[1])) = \begin{cases} \mathbb{R} & k=0 \\ 0 & k \neq 0 \end{cases}$$

cochain complex of vector spaces

\Rightarrow long exact sequence of vector balls

deg

-2 -1 0

$$\rightarrow S^*(V_x) \otimes \Lambda^2(V_x^\vee[1]) \xrightarrow{d} S^*(V_x) \otimes V_x^\vee[1] \xrightarrow{d} S^*(V_x) \xrightarrow{\epsilon} \mathbb{R}$$

exact
sequence
of vector
spaces

\Rightarrow exact sequence of vector bundles

final bar.

$$\rightarrow S^*(V^\vee) \otimes \wedge^*(V^\vee[1]) \rightarrow S^*(V^\vee) \otimes V^\vee[1] \rightarrow S^*(V^\vee) \rightarrow \underline{\mathbb{R}}$$



$$\rightarrow \Gamma(\quad) \rightarrow \Gamma(\quad) \rightarrow \Gamma(\quad) \xrightarrow{\xi} C^\infty(X)$$

$\underbrace{\Gamma(S^*(V^\vee) \otimes \wedge^*(V^\vee[1]))}$

Δ This exactness when passing to sections holds in the smooth setting!

So it suffices to show:

$$H^k(\underbrace{S^*(V_x^\vee) \otimes \wedge^*(V_x^\vee[1])}_{K(V)}) = \begin{cases} \mathbb{R} & k=0 \\ 0 & k \neq 0 \end{cases}$$

chain homotopy H_1 between \circ and a map \oint
to be determined.

$H : K(V) \rightarrow K(V^*)$ of degree -1 , a graded derivation,
 $S^k(V) \otimes \wedge^l(V^*[i])$ hence determined on generators

$$\xi \in V \quad H(\xi \otimes 1) := 1 \otimes \xi[i]$$

$$H(1 \otimes \xi[i]) := 0$$

$$\begin{aligned} \bar{\Phi} &:= dH + Hd \\ &= [d, H] \end{aligned} \quad \text{calculate, using the fact that}$$

graded commutator; hence since d & H are graded derivations
 $(\text{ob deg } + 1 \text{ resp. } -1)$

$$\bar{\Phi} = [d, H] \text{ graded derivation of deg } 0.$$

$$\bar{\Phi} : S^k(V^*) \otimes \wedge^l(V^*[i]) \hookrightarrow$$

is multiplication by $(k+l)$.

$\Rightarrow \bar{\Phi}$ is chain map to 0 ; hence it induces

Same map on cohomology.

Hence cohomology is concentrated or $\ell=0, k=0$ □

Recall: Cor: V vector bdL, section α
 \downarrow
 X zero section z

$$\mathcal{O}(X) \xrightarrow{h} \mathcal{O}(X) \cong \left(\Gamma\left(\wedge^*(V^*[1])\right), \iota_\alpha \right)$$

$z^*, \mathcal{O}(V), \alpha^*$

$$\iota_\alpha : \wedge^k(V^*[1]) \rightarrow \wedge^{k-1}(V^*[1])$$

in particular:

$$\mathcal{O}\left(\underbrace{\text{dCrt}(S)}_{\text{derived critical locus}}\right) = \mathcal{O}\left(X \xrightarrow{h} \text{graph}(\text{d}S)\right) = \left(\underbrace{\Gamma(X, \wedge^*(TX[1]))}_{\text{d}S}, \iota \right)$$

$\text{d}S : X \rightarrow T^*X$

diff. graded
algebra of
functions on
 $\text{dCrt}(S)$.

If dS is transversal to zero section, then

$$H^k(\Gamma(\Lambda^*(TM[1])), \iota_{dS}) = \begin{cases} C^\infty(\text{Crit}(S)) & k=0 \\ 0 & k \neq 0 \end{cases}$$

originally, we defined $\text{Obs}^{\text{cl}} := C^\infty(\text{Crit}(S))$
space of classical fields

from the derived point of view

we redefine

$$\underbrace{\text{Obs}^{\text{cl}}}_{\text{classical observables}} := \mathcal{O}(d\text{Crit}(S)) = (\Gamma(\Lambda^*(TX[1])), \iota_{dS}).$$

more explicitly: $\Gamma(\Lambda^*(TX[1])) \leftarrow$ a lot of structure,
 $\deg -3 \quad -2 \quad -1 \quad 0$ + (shifted) Lie algebra

$$\Gamma(\Lambda^3 TX) \xrightarrow{dS} \Gamma(\Lambda^2 TX) \xrightarrow{dS} \Gamma(TX) \xrightarrow{dS} C^\infty(X)$$

poly vector field

vector fields

Digression on symplectic manifolds
(classical mechanics)

Def: M mfd. A symplectic structure on M

is $\omega \in \Omega^2(M)$ s.t.

$$(i) \quad d\omega = 0$$

$$(ii) \quad T(TM) \xrightarrow{\cong} \Omega^1(M)$$

$$X \mapsto \iota_X \omega$$

$$(\iota_X \omega)(Y) = \omega(X, Y)$$

Rem: $\left\{ \begin{array}{l} \text{physical states} \\ \text{of a mechanical} \\ \text{system} \end{array} \right\}$ is typically a $\xrightarrow{\text{vector fields}}$ symplectic mfd.

Ex: a point particle moving in a Riem. mfd X
 $\left\{ \text{states} \right\} = \left\{ \text{geodesics in } X \right\} \cong TX \cong T^*_{\text{Riem. metric}} X$

for any mfd X (without any structure),
 there is a preferred symplectic structure
 on T^*X :

$$\omega = d\Theta \in \Omega^2(T^*X)$$

$$\Theta((x, \xi), v) := \xi(\pi_*(v)) \in \mathbb{R}$$

$$x \in T_x^*X \quad T_{(x, \xi)}(T^*X)$$

$$\pi: T^*X \rightarrow X \quad \text{differential } \pi_*: T_{(x, \xi)} T^*X \rightarrow T_x X$$

We are interested in the time development of
 $\phi \in M$ (symplectic mfd).

Def.: given $f \in C^\infty(M)$, the associated
Hamiltonian vector field X_f
 is determined by $\iota_{X_f} \omega = -df$

$$\text{Ex: } M = T^* \mathbb{R} = \overset{\text{base}}{\mathbb{R}} \times \overset{\text{fiber}}{\mathbb{R}}$$

coordinate functions: q, p
 "position" "momentum"

homework: i) $\Theta = pdq$ and hence $\omega = d\Theta = dp \wedge dq$
 ii) $X_q = -\frac{\partial}{\partial p}, X_p = \frac{\partial}{\partial q}$.

Think of functions on a symplectic mfd M
 as "observables".

Def: Poisson bracket: $\{, \}: C^\infty(M) \times C^\infty(M) \rightarrow C^\infty(M)$
 M symplectic

$$\{f, g\} := X_f g = dg(X_f) = -(L_{X_g} \omega)(X_f) = -\omega(X_g, X_f) = \omega(X_f, X_g)$$

Lem: The Poisson bracket gives $C^\infty(M)$ the structure of a Lie algebra, i.e.

- (skew-symmetry) $\{f, g\} = -\{g, f\}$
- (Jacobi identity): $\{\{f, g\}, h\} + \{\{g, h\}, f\} + \{\{h, f\}, g\} = 0$

$$\text{i.e. } \{f, -\} \text{ is a derivation w.r.t. to the Lie bracket.}$$
$$\{f, \{g, h\}\} \xrightarrow{\text{def}} \{\{f, g\}, h\} + \{g, \{f, h\}\}$$

Pf: homework

Lem: $C^\infty(M) \longrightarrow T(TM)$

given by $\{, \}$ $\xrightarrow{\text{Lie alg.}}$ $f \longmapsto X_f$ $\xrightarrow{\text{Lie bracket of vector fields}}$ is a Lie algebra homomorphism,

$$\text{i.e. } [x_f, x_g] = X_{\{f, g\}}$$