## An Eigenvalue-Based Estimate of Reaction Zone

## Thicknesses in Gas Phase Detonations

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## Motivation

- Detailed kinetics models are widely used in detonation simulations.
- The finest length scale predicted by such models is usually not clarified and often not resolved.
- Tuning computational results to match experiments without first harmonizing with underlying mathematics renders predictions unreliable.


## Model: Reactive Euler Equations

- one-dimensional
- steady
- inviscid
- detailed Arrhenius kinetics
- calorically imperfect ideal gas mixture


## Model: Reactive Euler Equations

$$
\begin{aligned}
\rho u & =\rho_{o} D \\
\rho u^{2}+p & =\rho_{o} D^{2}+p_{o} \\
e+\frac{u^{2}}{2}+\frac{p}{\rho} & =e_{o}+\frac{D^{2}}{2}+\frac{p_{o}}{\rho_{o}} \\
\frac{d Y_{i}}{d x} & =f_{i} \equiv \frac{\dot{\omega}_{i} M_{i}}{\rho_{o} D}
\end{aligned}
$$

Supplemented by state equations and the law of mass action.

## Reduced Model

Algebraic reductions lead to a final form of

$$
\frac{d Y_{i}}{d x}=f_{i}\left(Y_{1}, \ldots, Y_{N-L}\right)
$$

with

- $N$ : number of molecular species
- $L$ : number of atomic elements


## Eigenvalue Analysis of Local Length Scales

Local behavior is modelled by

$$
\frac{d \mathbf{Y}}{d x}=\mathbf{J} \cdot\left(\mathbf{Y}-\mathbf{Y}^{*}\right)+\mathbf{b}, \quad \mathbf{Y}\left(x^{*}\right)=\mathbf{Y}^{*}
$$

whose solution has the form

$$
\mathbf{Y}(x)=\mathbf{Y}^{*}+\left(\mathbf{P} \cdot e^{\boldsymbol{\Lambda}\left(x-x^{*}\right)} \cdot \mathbf{P}^{-1}-\mathbf{I}\right) \cdot \mathbf{J}^{-1} \cdot \mathbf{b}
$$

Here $\boldsymbol{\Lambda}$ has eigenvalues $\lambda_{i}$ of Jacobian $\mathbf{J}$ in its diagonal.
The length scales are given by

$$
\ell_{i}(x)=\frac{1}{\left|\lambda_{i}(x)\right|}
$$

## Computational Methods

- A standard ODE solver (DLSODE) was used to integrate the equations.
- Standard IMSL subroutines were used to evaluate the local Jacobians and eigenvalues at every step.
- The Chemkin software package was used to evaluate kinetic rates and thermodynamic properties.
- Computation time was typically two minutes on a 900 MHz Sun Blade 1000.


## Physical System

- Hydrogen-air detonation: $2 \mathrm{H}_{2}+\mathrm{O}_{2}+3.76 \mathrm{~N}_{2}$.
- $N=9$ molecular species, $L=3$ atomic elements, $J=19$ reversible reactions.
- $p_{o}=1 \mathrm{~atm}$.
- $T_{o}=298 \mathrm{~K}$.
- Identical to system studied by both Shepherd (1986) and Mikolaitis (1987).


## Mole Fractions versus Distance



- significant evolution at fine length scales $x<10^{-3} \mathrm{~cm}$.


## Temperature Profile



- Temperature flat in the post-shock induction zone $0<x<2.6 \times 10^{-2} \mathrm{~cm}$.
- Thermal explosion followed by relaxation to

Eigenvalue Analysis: Length Scale Evolution


- Finest length scale: $2.3 \times 10^{-5} \mathrm{~cm}$.
- Coarsest length scale $3.0 \times 10^{1} \mathrm{~cm}$.


## Influence of Initial Pressure



- Induction zone length and finest length scale are sensitive to initial pressure.
- Finest lenath scale three orders of maanitude


## Verification: Comparison with Mikolaitis



- Lagrangian calculation allows direct comparison with Mikolaitis' results.


## Grid Convergence



- Finest length scale must be resolved to converge at proper order.
- Results are converging at proper order for first


## Numerical Stability



- Discretizations finer than finest physical length scale are numerically stable.
- Discretizations coarser than finest physical

Examination of Recently Published Results

| Ref. | $\ell_{\text {ind }}(\mathrm{cm})$ | $\ell_{f}(\mathrm{~cm})$ | $\Delta x(\mathrm{~cm})$ |
| :---: | :---: | :--- | :--- |
| Oran, et al, 1998 | $1.47 \times 10^{-1}$ | $2.17 \times 10^{-4}$ | $3.88 \times 10^{-3}$ |
| Shepherd, et al., 2003 | $1.46 \times 10^{-1}$ | $2.16 \times 10^{-4}$ | AMR |
| Hayashi, et al., 2002 | $1.50 \times 10^{-2}$ | $1.23 \times 10^{-5}$ | $5.00 \times 10^{-4}$ |
| Powers, et al., 2001 | $1.54 \times 10^{-2}$ | $2.76 \times 10^{-5}$ | $8.14 \times 10^{-5}$ |
| Fedkiw, et al., 1997 | $1.54 \times 10^{-2}$ | $2.76 \times 10^{-5}$ | $3.00 \times 10^{-2}$ |
| Ebrahimi and Merkle, 2002 | $5.30 \times 10^{-3}$ | $7.48 \times 10^{-6}$ | $1.00 \times 10^{-2}$ |
| Sislian, et al., 1998 | $1.38 \times 10^{-1}$ | $2.23 \times 10^{-4}$ | $1.00 \times 10^{0}$ |

All are under-resolved, some severely.

## Conclusions

- Detonation calculations are often under-resolved, by as much as four orders of magnitude.
- Equilibrium properties are insensitive to resolution, while transient phenomena can be sensitive.
- Sensitivity of results to resolution is not known a priori.
- Numerical viscosity stabilizes instabilities.
- For a repeatable scientific calculation of detonation, the finest physical scales must be resolved.


## Moral

You either do detailed kinetics with the proper resolution,

or

you are fooling yourself and others, in which case you should stick with reduced kinetics!

