

Perspectives from Poisson Geometry

Arlo Caine



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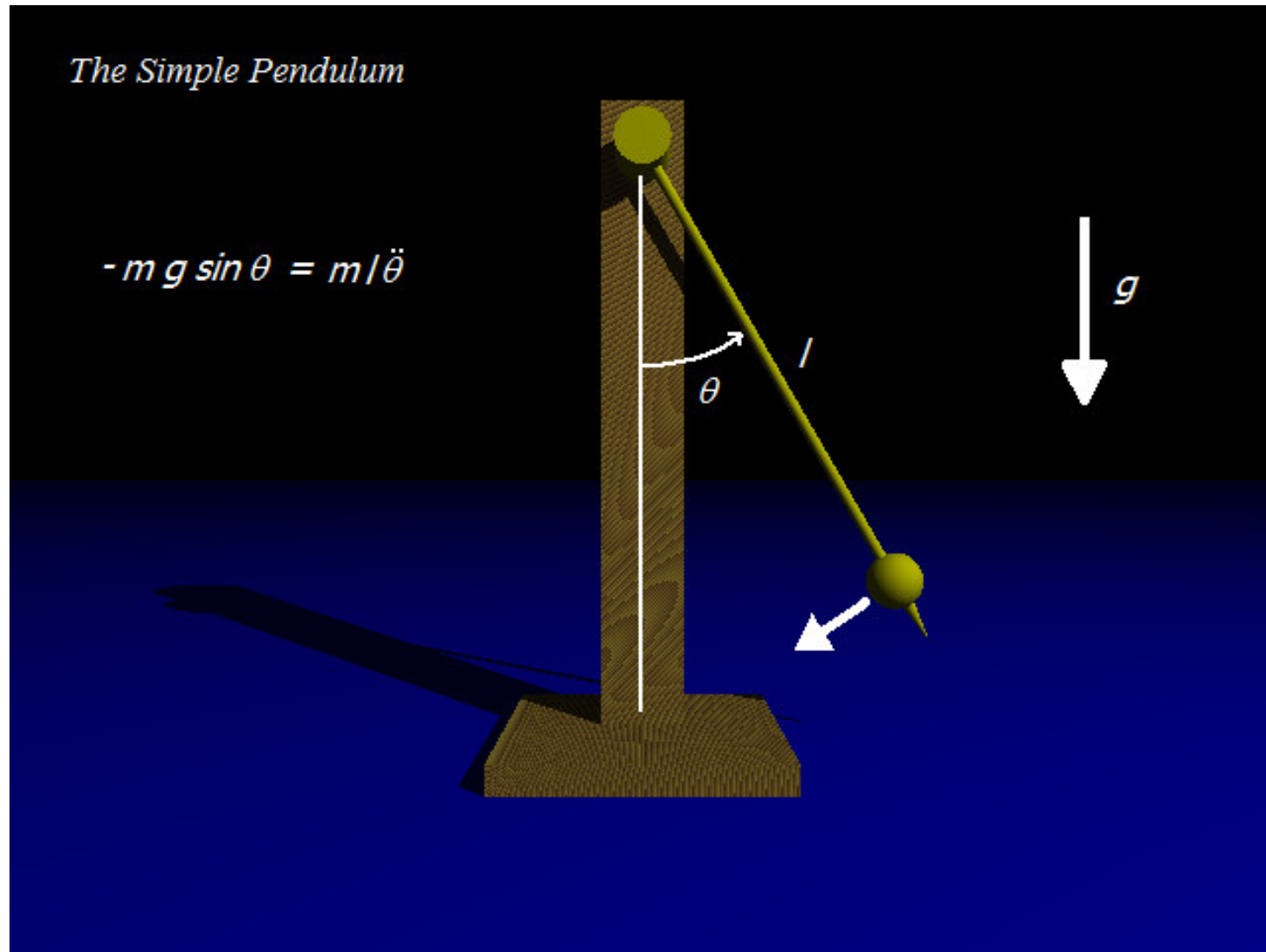
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Colloquium: University College Cork

29 February 2008



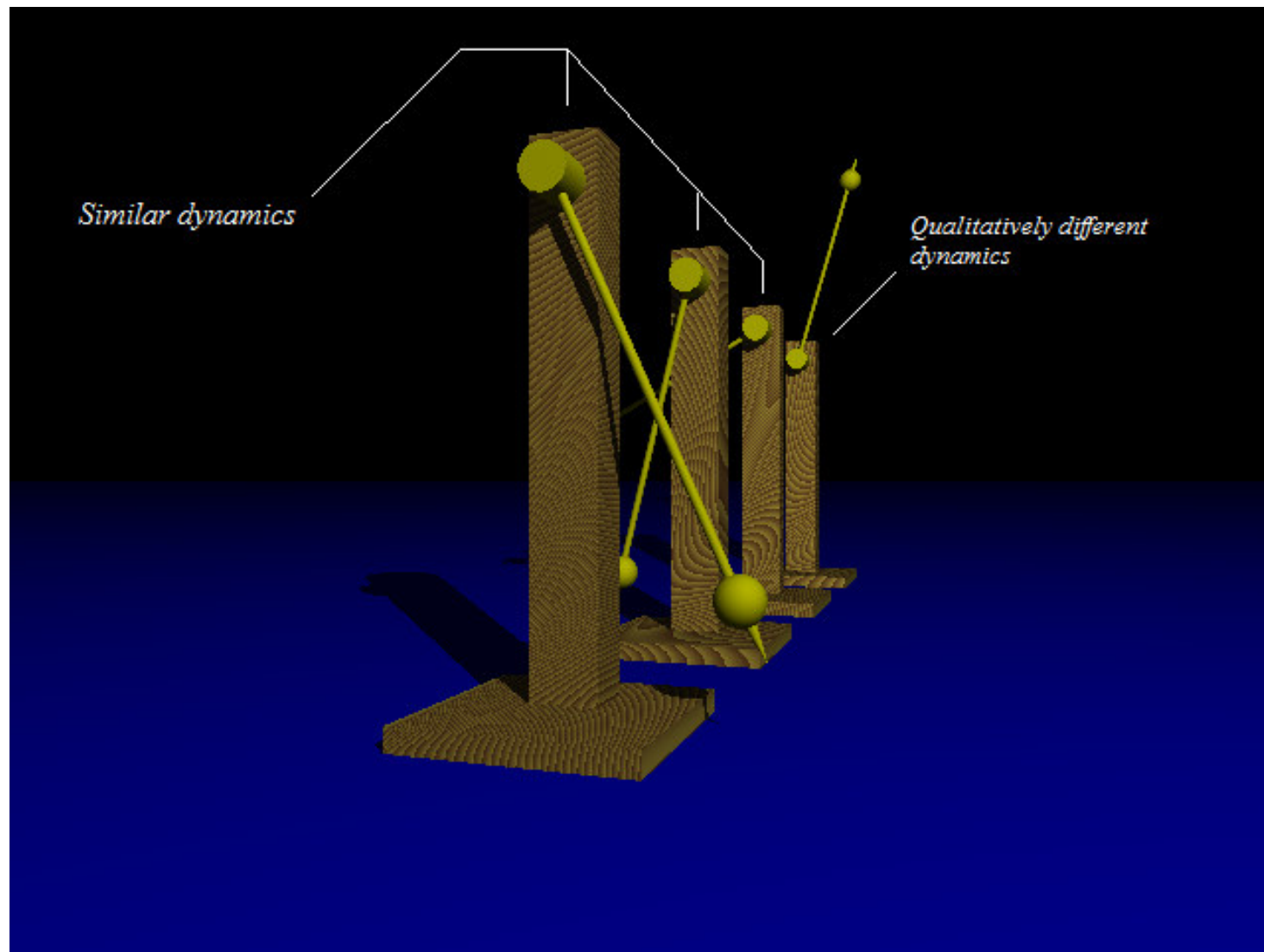
Mechanics of Motion



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Mechanics of Motion



Hamilton's Idea

Look at the E.O.M. in first order form

$$\begin{cases} \dot{v} &= -g \sin(q/\ell) \\ \dot{q} &= v \end{cases}$$

Rewrite in terms of linear momentum $p = mv$.

$$\begin{cases} \dot{p} &= -mg \sin(q/\ell) \\ \dot{q} &= p/m \end{cases}$$

Key Observation:

$$\begin{cases} \dot{p} &= -\partial H / \partial q \\ \dot{q} &= \partial H / \partial p \end{cases}$$

where H is the total **Total Energy**

$$H = \underbrace{\frac{1}{2m} p^2}_{\text{kinetic energy}} + \underbrace{(-mgl \cos(q/\ell) + mgl)}_{\text{potential energy}}.$$

$$q = \ell\theta, v = \dot{q} = \ell\dot{\theta}$$



www.ria.ie/



Conserved Quantities

H is constant along the flow. (Conservation of Energy)

$$\frac{d}{dt}H = \frac{\partial H}{\partial q}\dot{q} + \frac{\partial H}{\partial p}\dot{p}$$



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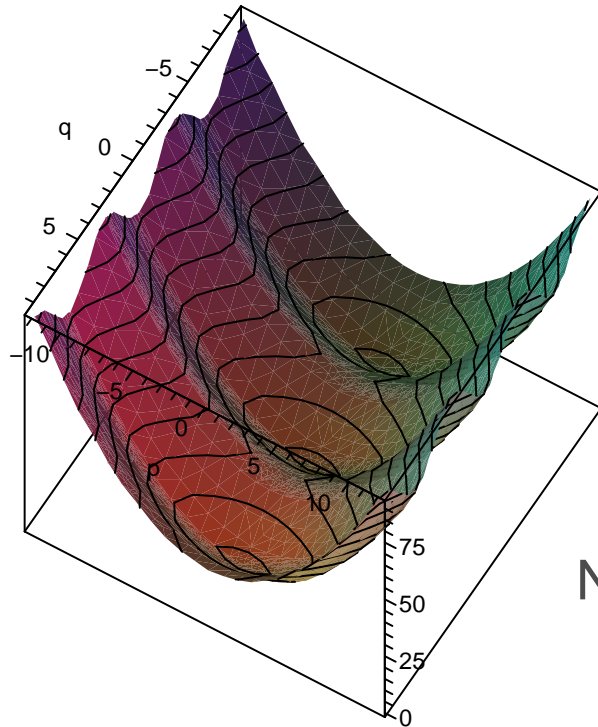
$$\frac{d}{dt}H = \frac{\partial H}{\partial q}\dot{q} + \frac{\partial H}{\partial p}\dot{p} \stackrel{HE}{=} \frac{\partial H}{\partial q}\frac{\partial H}{\partial p} + \frac{\partial H}{\partial p}\left(-\frac{\partial H}{\partial q}\right) = 0$$



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- 1 Dimensional Problem:
Qualitative view of dynamics obtained by analyzing the level sets of H .
- Higher Dimensions:

Need **independent** conserved quantities.

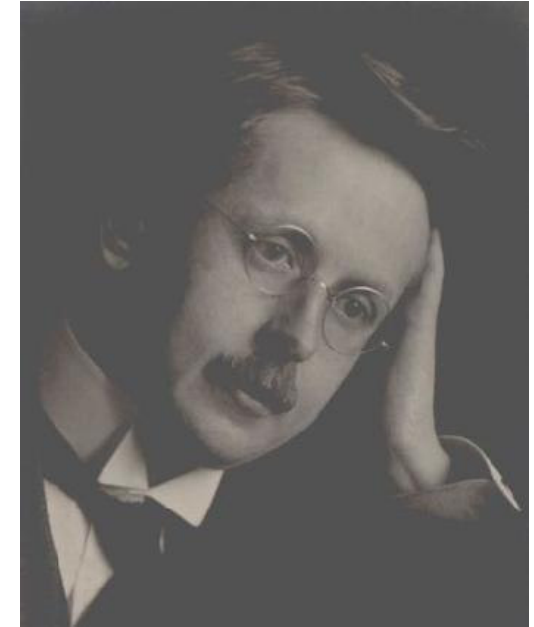


Weyl's Observations

- Velocities are *tangent vectors*.
- Physics \Rightarrow Momentum is a *1-form*.
- Why? Pair with velocity \rightarrow kinetic energy.
- View H as a function $H: T^*Q \rightarrow \mathbb{R}$.
- Hamilton's E.O.M. \Leftrightarrow equality of 1-forms

$$\omega(X, \cdot) = dH(\cdot)$$

where $X = \dot{q} \frac{\partial}{\partial q} + \dot{p} \frac{\partial}{\partial p}$ and $\omega = dp \wedge dq$.



www.weylmann.com

$$\begin{pmatrix} \dot{q} & \dot{p} \end{pmatrix} \underbrace{\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}}_J = \begin{pmatrix} \frac{\partial H}{\partial q} & \frac{\partial H}{\partial p} \end{pmatrix} \rightsquigarrow J^2 = -1$$

complex \Leftrightarrow symplectic



Poisson's Viewpoint

- Hamilton's E.O.M are equivalent to:

$$X = \pi(\cdot, dH)$$

where $\pi = \frac{\partial}{\partial q} \wedge \frac{\partial}{\partial p}$ (bivector on T^*Q).



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$$\{F_1, F_2\} := \pi(dF_1, dF_2).$$



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Key Point:

$$\{\{F_1, F_2\}, H\} + \{\{F_2, H\}, F_1\} + \{\{H, F_1\}, F_2\} = 0.$$

Thus, F_1 and F_2 conserved $\Rightarrow \{F_1, F_2\}$ is conserved.



Geometric Abstraction

Geometric structures on M .

- **Definition 1:** A *symplectic structure* on M is a 2-form ω on M such that:
 1. ω is non-degenerate on $T_p M$ for each $p \in M$,
 2. $d\omega = 0$.



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 1. $\{f, g\} = -\{g, f\}$,
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2 & 3 are equivalent: $\{f, g\} := \pi(df, dg)$

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Symplectic Foliation

Anchor Map:

$$\begin{aligned}\pi^\# : T^*M &\rightarrow TM \\ \alpha &\mapsto \pi(\cdot, \alpha)\end{aligned}$$

- $\text{Image}(\pi^\#|_p) \subset T_pM$ is even dimensional for each $p \in M$.
- $\text{Rank}(\pi^\#|_p)$ is not necessarily constant in p .
- $\text{Image}(\pi^\#)$ defines a general distribution Δ in TM .
- π is tangent to $\text{Image}(\pi^\#)$ and non-degenerate there.
- $[\pi, \pi]_{Sch} = 0 \stackrel{SS}{\Rightarrow} \Delta$ is integrable and π^{-1} is symplectic on each leaf.

Sussmann-Stefan Frobenius Theorem



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Poisson manifolds are foliated by symplectic submanifolds

Sussmann-Stefan Frobenius Theorem



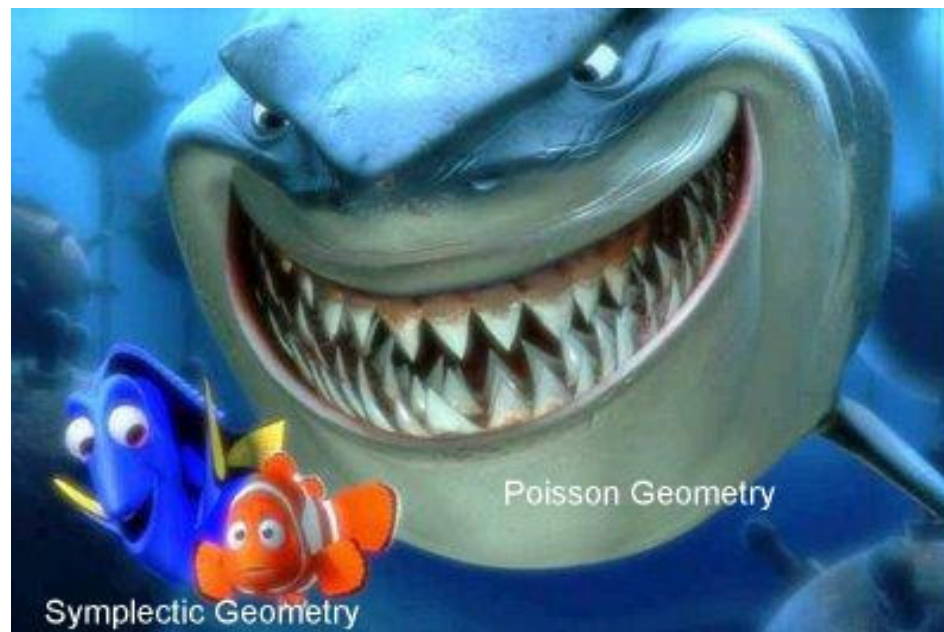
General Examples

- World's most uninteresting example: $\pi = 0$ on any M .
- (M^{2k}, ω) symplectic $\Rightarrow \pi = \omega^{-1}$ is a Poisson structure on M and Δ has exactly one symplectic leaf.
- (M^2, ω) symplectic $\pi = f\omega^{-1}$ is Poisson for each $f \in C^\infty(M)$. Δ has zero dimensional leaves at the zeros of f .



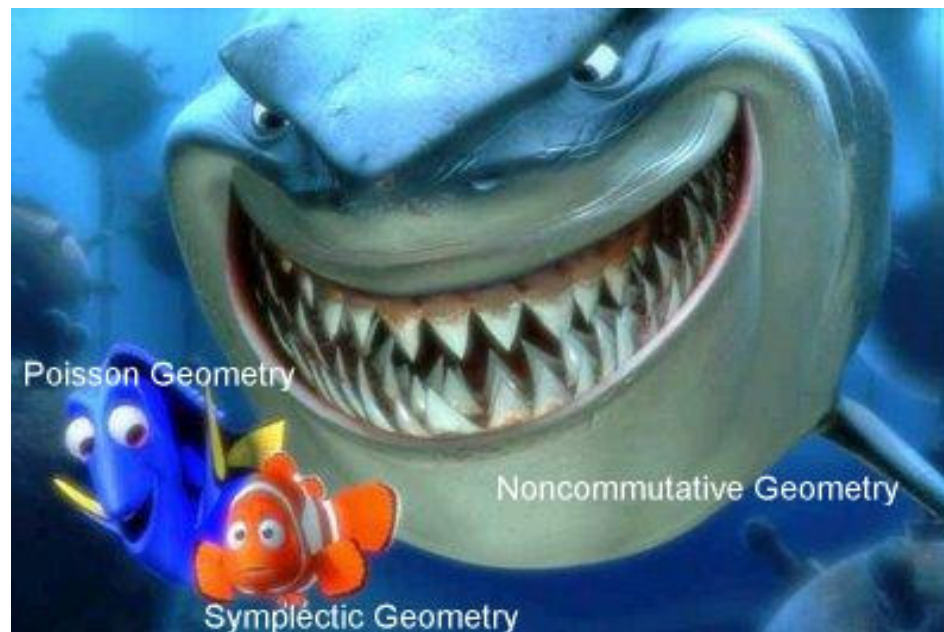
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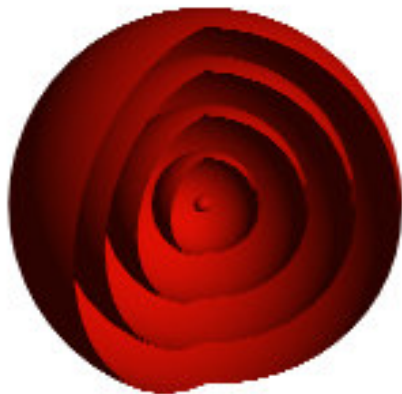
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Poisson Structures on \mathbb{R}^3

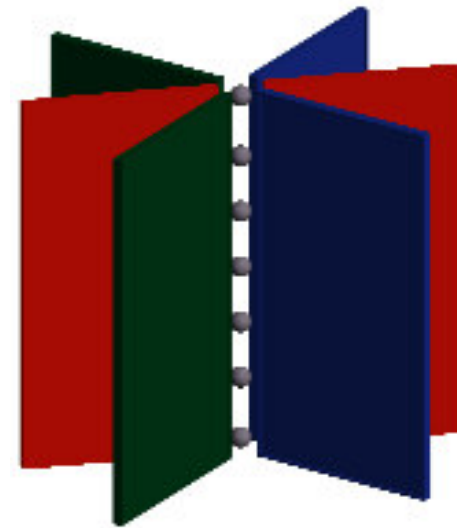
$\{f, g\} := (\nabla f \times \vec{F}) \cdot \nabla g$ where $\vec{F} \cdot (\nabla \times \vec{F}) = 0$.



$$\vec{F} = \nabla\left(\frac{1}{2}(x^2 + y^2 + z^2)\right)$$



$$\vec{F} = \nabla\left(\frac{1}{2}(x^2 + y^2)\right)$$



$$\vec{F} = -y\vec{i} + x\vec{j}$$

$$\vec{F} = P\vec{i} + Q\vec{j} + R\vec{k} \rightsquigarrow \pi = P\partial_y \wedge \partial_z + Q\partial_z \wedge \partial_x + R\partial_x \wedge \partial_y$$

$$C^\infty(\mathbb{R}^3)/\mathbb{R} \hookrightarrow \{\text{Poisson Structures on } \mathbb{R}^3\}$$

Lots of variety!

Non-gradient examples \leftrightarrow Triply orthogonal coordinate systems



Lie Poisson Structures

Recall $\{f, g\} = \pi(df, dg)$.

- $\{\cdot, \cdot\}$ is a derivation in each slot.
- Enough to specify its action on *coordinate functions*.



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Let \mathfrak{g} be a finite dimensional real Lie algebra.

- View elements of \mathfrak{g} as linear functions on \mathfrak{g}^* .
- Basis X_1, \dots, X_n of \mathfrak{g} determines linear coordinates on \mathfrak{g}^* .
- Define $\{\cdot, \cdot\}$ on \mathfrak{g}^* by

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Theorem: (Weinstein) Symplectic leaves are coadjoint orbits.



Poisson Lie Groups

New feature!

- **Definition:** A Poisson map $f : (M, \pi_M) \rightarrow (N, \pi_N)$ is a smooth map such that $f_*\pi_M = \pi_N$.
- **Definition:** A Poisson Lie group is a Lie group U with a Poisson structure π_U such that the multiplication map $m : U \times U \rightarrow U$ is a Poisson map, i.e. $m_*(\pi_U \oplus \pi_U) = \pi_U$.

Notes:

- $\pi_U = 0$ is a (boring) Poisson Lie group structure on U .
- $(\cdot)^{-1} : U \rightarrow U$ is Poisson if and only if $\pi_U = 0$.

Interesting Examples?



Standard Example

Theorem:(Drinfel'd) $\pi_U = \Lambda^r - \Lambda^\ell$ is a Poisson Lie group structure if and only if $[\Lambda, \Lambda] \in \wedge^3 \mathfrak{u}$ is Ad-invariant.



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Theorem: $[\mathcal{H}, \mathcal{H}]_{Sch}$ is an ad-invariant element of $\wedge^3 \mathfrak{u}$.

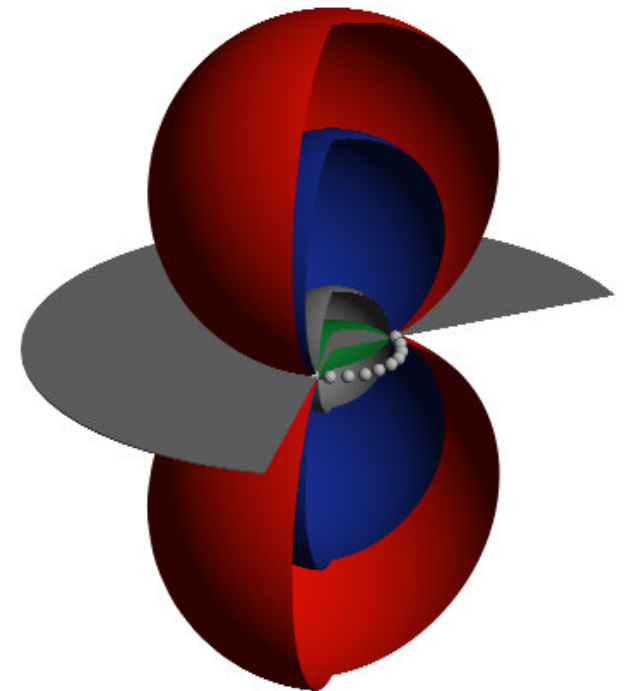
Many equivalent versions, observed by many authors.



On $SU(2)$

$$\begin{pmatrix} a & b \\ -\bar{b} & \bar{a} \end{pmatrix} \mapsto \pi_U = (1 - |a|^2 - |b|^2)X \wedge Y + 2\operatorname{Im}(a\bar{b})Y \wedge H - 2\operatorname{Re}(a\bar{b})H \wedge X.$$

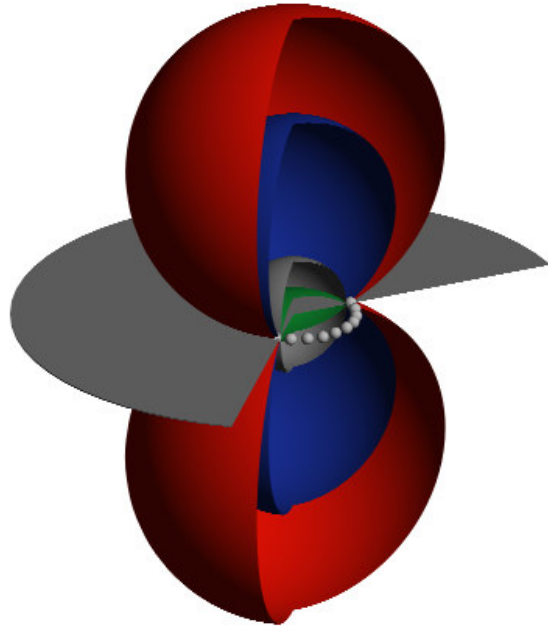
- $\Lambda^2 TSU(2) \simeq SU(2) \times \Lambda^2 \mathfrak{su}(2)$
- $\mathfrak{su}(2) = \operatorname{span}_{\mathbb{R}}\{H, X, Y\}$
(i -Pauli spin matrices)
- π_U vanishes on $T = \{b = 0\}$.
- Δ has zero dimensional leaves along T .
- 2-d leaves are disks whose boundary meets T .



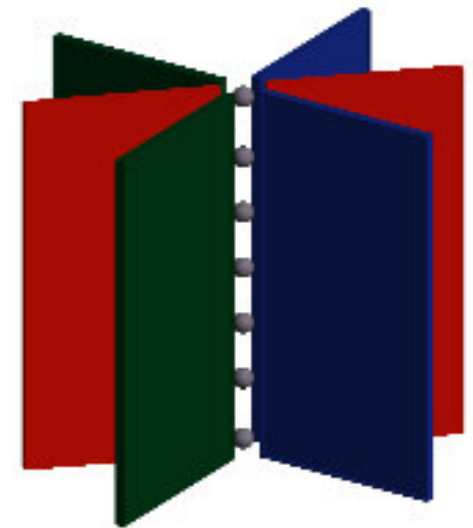
Stereographic picture in S^3



From another viewpoint



$$\infty \leftrightarrow \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$$



$$\infty \leftrightarrow \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Generates a non-gradient example of 3-dimensional Poisson structure!



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Example:

- $SU(2)/T \simeq \mathbb{C}P^1$
- The two symplectic leaves give a cell decomposition of $\mathbb{C}P^1$.



Bruhat Poisson Structure on U

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} 1 & \frac{a}{c} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \frac{c}{1} & 0 \\ 0 & \frac{1}{c} \end{pmatrix} \begin{pmatrix} 1 & \frac{d}{c} \\ 0 & 1 \end{pmatrix}$$

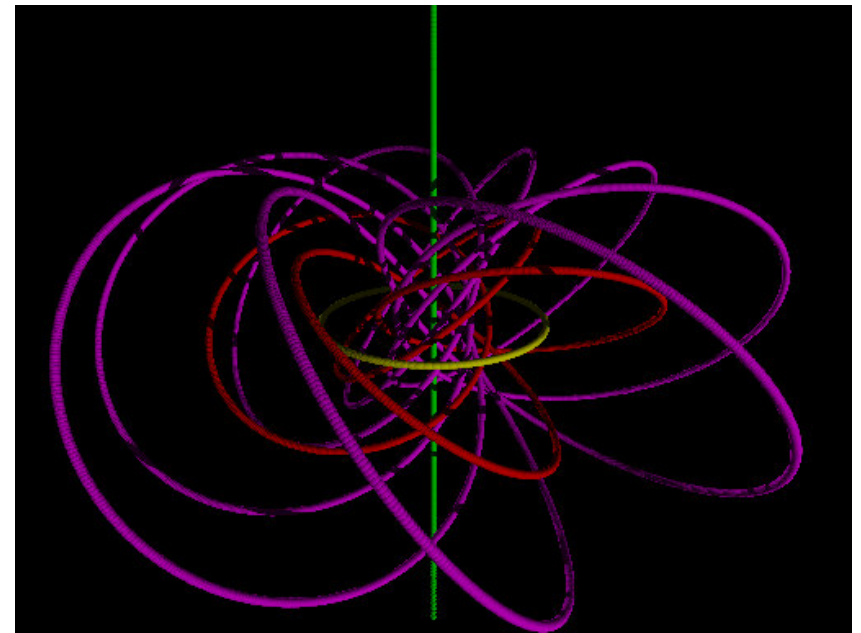
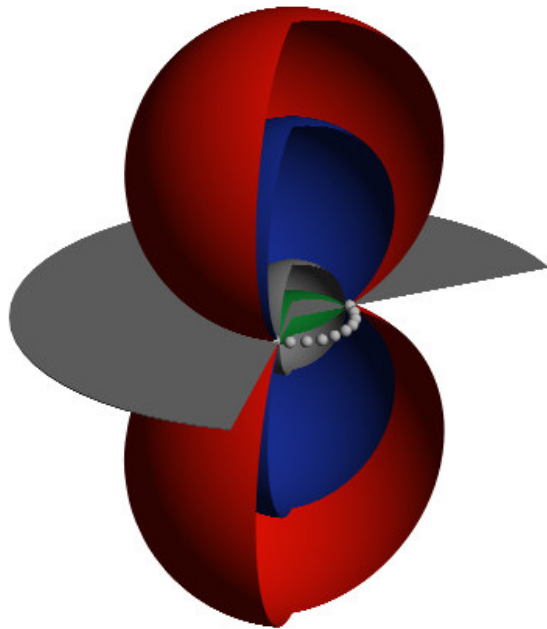
$$G = \coprod_{w \in W} N^+ w H N^+$$



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Birkhoff Decomposition of U

$$\begin{pmatrix} 1 & 0 & 0 \\ \frac{a_{21}}{\sigma_1} & 1 & 0 \\ \frac{a_{31}}{\sigma_1} & \frac{M_{23}}{\sigma_2} & 1 \end{pmatrix} \begin{pmatrix} \frac{\sigma_1}{1} & 0 & 0 \\ 0 & \frac{\sigma_2}{\sigma_1} & 0 \\ 0 & 0 & \frac{1}{\sigma_2} \end{pmatrix} \begin{pmatrix} 1 & \frac{a_{12}}{\sigma_1} & \frac{a_{13}}{\sigma_1} \\ 0 & 1 & \frac{M_{32}}{\sigma_2} \\ 0 & 0 & 1 \end{pmatrix}$$

- $A = (a_{ij}) \in SL(3, \mathbb{C})$ with non-trivial principal minors σ_1, σ_2 .
- Gaussian elimination \Rightarrow factorization exists and is unique for $\sigma_1\sigma_2 \neq 0$.
- Introduction of permutation factor necessary for higher codimension components (e.g. $\sigma_1 = 0$.)
- Birkhoff Decomposition

$$G = \coprod_{w \in W} \Sigma_w^G \text{ where } \Sigma_w^G = N^- w H N^+$$

- Induced decomposition $U = \coprod_{w \in W} (U \cap \Sigma_w^G)$.



Birkhoff Poisson Structure

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- **Key Point:** Longest element of the Weyl group interchanges \mathfrak{n}^+ and \mathfrak{n}^- , sending \mathcal{H} to $-\mathcal{H}$.



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- Π_U arises by applying the Evens-Lu construction to generate a $(U \times U, \pi_U \oplus \pi_U)$ -homogeneous Poisson structure on U (viewed as a symmetric space).

Theorem:(C,Pickrell) The Poisson structures π_U and Π_U are Poisson equivalent by translation by the longest element of the Weyl group followed by inversion.

- **Key Point:** Longest element of the Weyl group interchanges \mathfrak{n}^+ and \mathfrak{n}^- , sending \mathcal{H} to $-\mathcal{H}$.

Final Comment:

- The first theorem can be translated in to the infinite dimensional setting of the loop group of U , but the second result is finite dimensional.

