

ASSIGNMENT 4 SOLUTIONS

(a)  $V_{bi} = \frac{kT}{q} \ln\left(\frac{N_A N_D}{n_c^2}\right) = \boxed{1.22 \text{ V}}$

PROBLEM 1 SOLUTION

$W = \left[ \frac{2\epsilon_s}{q} \left( \frac{1}{N_A} + \frac{1}{N_D} \right) V_{bi} \right]^{1/2} \approx \boxed{440 \text{ nm}}$

Since  $\frac{N_A}{N_D} = 10$ ,  $x_n = \frac{1}{1+10} * W \approx \boxed{40 \text{ nm}}$

$x_p = \frac{10}{1+10} * W \approx \boxed{400 \text{ nm}}$

$x_p \sim 10 x_n$

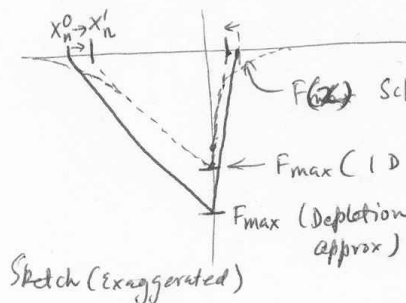
$F_{max} = \frac{2V_{bi}}{W} \approx \boxed{55.5 \text{ kV/cm}}$

$\ll F_{breakdown} = 400 \text{ kV/cm}$

$C_{depl} = \frac{\epsilon_s}{W} \approx \boxed{26 \text{ nF/cm}^2}$

Sketches & Plots: Next pages by QIN.

(b) Plots - next pages.



$W_{depl}(Gummel \text{ correction}) < W_{depl}(Depl. approx)$

$F_{max}(\text{''}) < F_{max}(\text{''})$

Schrodinger/Poisson soln.

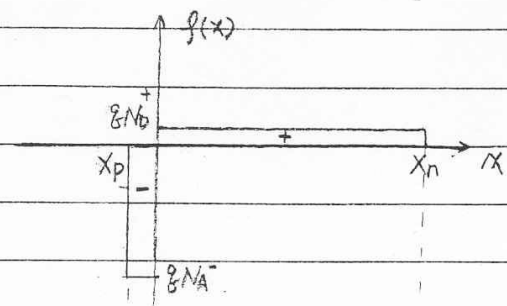
$F_{max}$  (1D Poisson, with Gummel correction).

$F_{max}$  (Depletion approx)

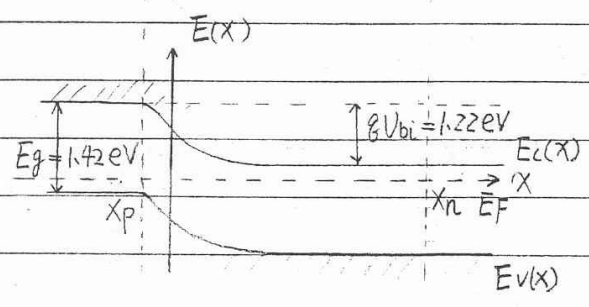
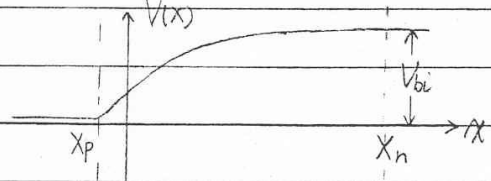
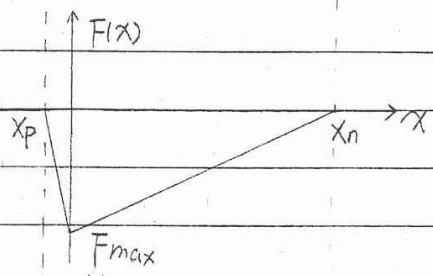
BUT:

$V_{bi}$  remains the SAME!!

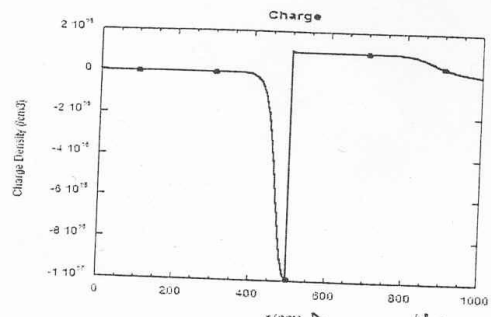
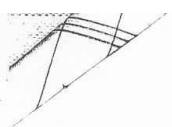
2)



Plot - by Qin.

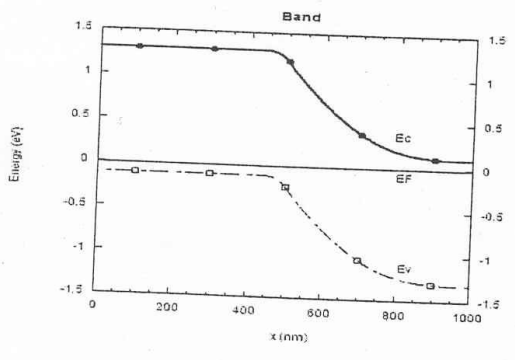
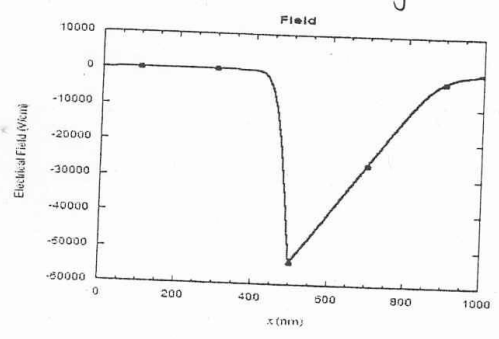


3



Plot - by  $\Phi_{in}$

Use the junction as  $x=0$ .



PROBLEM 2 SOLUTION

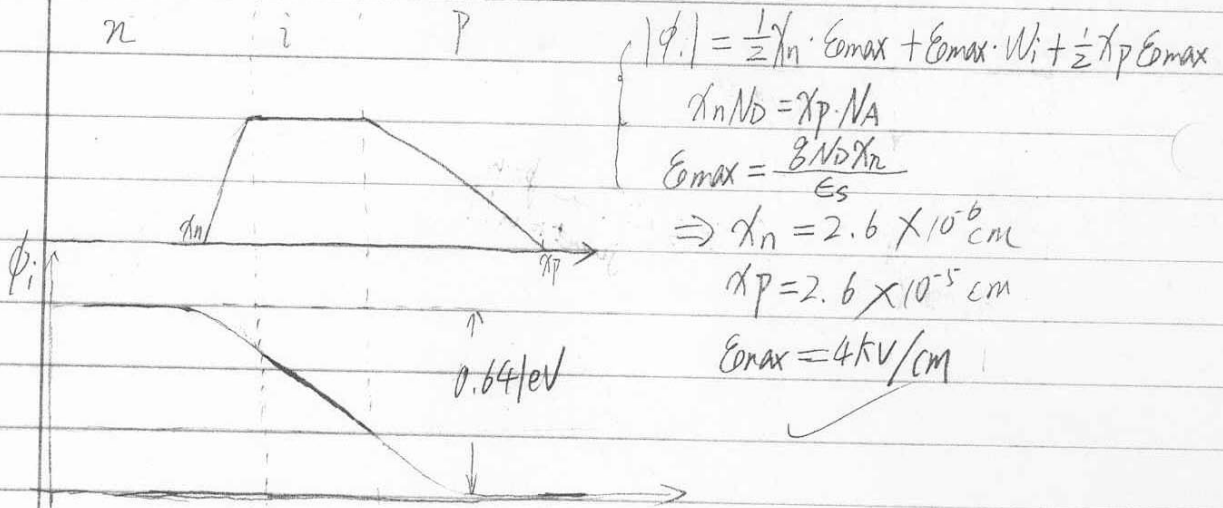
MKC 4.5 a) From the figure, we can consider the case as abrupt junction.

Consider in n-region  $N_D = 1 \times 10^{16}$ , in p-region  $N_A = 0.5 \times 10^{15}$

$$\phi_i = kT \ln \left( \frac{N_A N_D}{n_i^2} \right)$$

$$= 0.026 \cdot \ln \left( \frac{10^{16} \times 10^{15}}{(1.4 \times 10^{10})^2} \right) = 0.64 \text{ eV}$$

$\therefore \phi_i = 0.64 \text{ V}$



b) In p-n junction.

$\phi_i = 0.64 \text{ V}$

$$|\phi_i| = \frac{1}{2} E_{max} (x_n + x_p)$$

$$N_D x_n = N_A x_p$$

$$E_{max} = \frac{q N_D x_n}{\epsilon_s}$$

$\Rightarrow E_{max} = 1.3 \times 10^5 \text{ V/cm}$

$\therefore E_{max \text{ p-n}} > E_{max \text{ p-i-n}}$

### PROBLEM 3 SOLUTION

problem 4. a) i)  $E_f - E_i = kT \ln \left( \frac{N_c}{n_i} \right) = 0.026 \times \ln \left( \frac{5 \times 10^{15}}{1.4 \times 10^{10}} \right) = 0.33 \text{ eV}$

ii)  $P = 10 \text{ W/cm}^2$  From fig 1.16 we can find the value  $\mu_p = 460 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$

$P = \frac{1}{84qP} \Rightarrow P \approx 1.5 \times 10^{15} \text{ cm}^{-3}$

$E_i - E_f = kT \ln \left( \frac{P}{n_i} \right) = 0.026 \times \ln \left( \frac{1.5 \times 10^{15}}{1.4 \times 10^{10}} \right) = 0.3 \text{ eV}$

b)  $\phi_{B0} = 0.85 \text{ eV}$

i)  $\phi_{B1} = \phi_{B0} - (E_f - E_i)$

$= \phi_{B0} - kT \ln \frac{N_c}{n_0} = 0.85 - 0.026 \cdot \ln \frac{3.2 \times 10^7}{5 \times 10^{15}}$

$= 0.622 \text{ eV}$

$\phi_1 = 0.622 \text{ Volt}$

(ii)  $\phi_{B2} = \phi_{B0} - \phi_{B1} = 6.125 - 4.05 = 2.075$

(For the work function of  $\text{Pt}$ , I can find different value from  $5.9 \text{ eV} \sim 6.35 \text{ eV}$ . I choose the average number  $\frac{5.9 + 6.35}{2} = 6.125$

The  $0.85 \text{ eV}$  barrier is not consistent with idealized Schottky theory. The reason is Fermi level pinning, which makes the  $\phi_{B0}$  almost independent on the kind of metal  $\phi_{B0} \approx \frac{2}{3} E_g = 0.75$

c. when complete depleted,  $x_d = 2.5 \mu\text{m}$

$$E_{\text{max}} = \frac{qN_A x_d}{\epsilon_s} = \frac{1.62 \times 10^{-19} \times 5 \times 10^{15} \times 2.5 \times 10^{-4}}{1.7 \times 8.85 \times 10^{-14}}$$

$$= 1.93 \times 10^5 \text{ V/cm}$$

∴ It is possible to deplete the whole n-layer without reaching the breakdown voltage of  $3 \times 10^5 \text{ V/cm}$ .

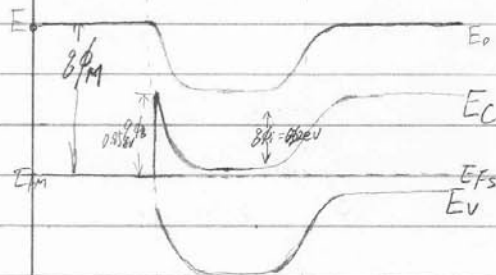
when completely depleted,  $x_d = 2.5 \mu\text{m}$

$$\frac{1}{2} E_{\text{max}} x_d = \frac{qN_A x_d^2}{2\epsilon_s} = \phi - V_a$$

$$\Rightarrow V_a = 23.5 \text{ V}$$

reversed bias.

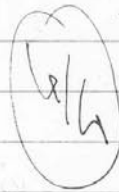
d. M n<sup>-</sup> P



$$q\phi_i = kT \ln \left( \frac{N_D N_A}{n_i^2} \right)$$

$$= 0.026 \times \ln \left( \frac{5 \times 10^{15} \times 1.5 \times 10^{15}}{(1.4 \times 10^{10})^2} \right)$$

$$= 0.62 \text{ eV}$$



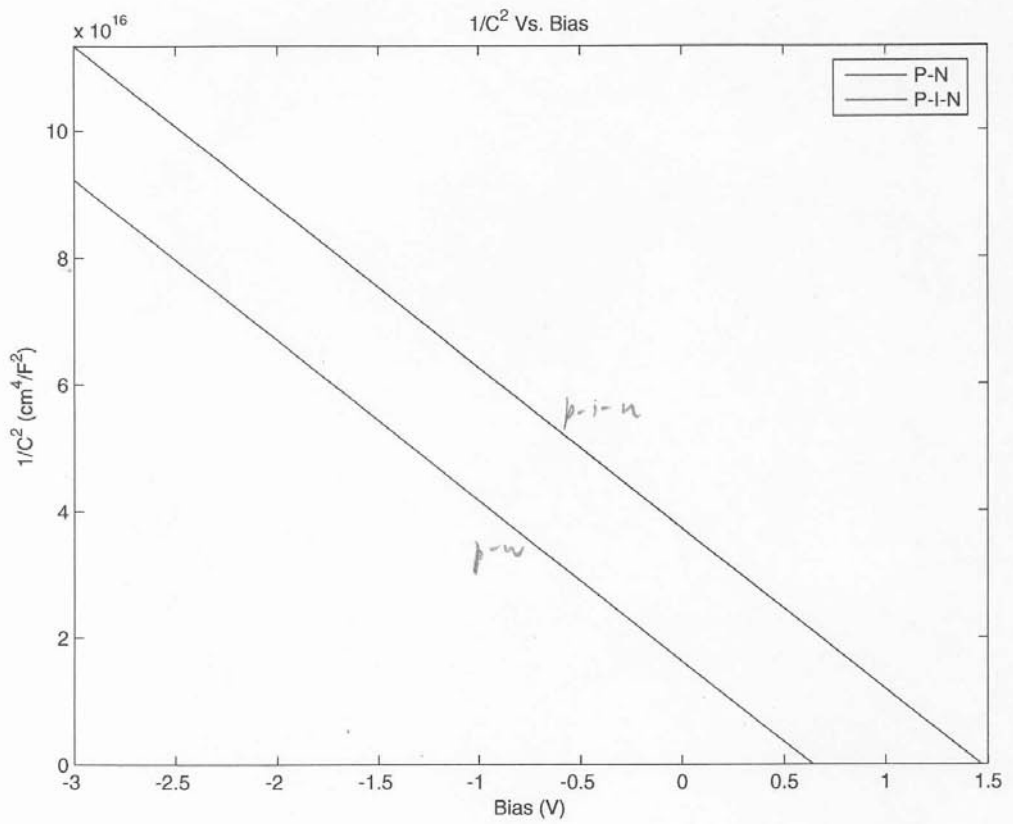


Fig 4

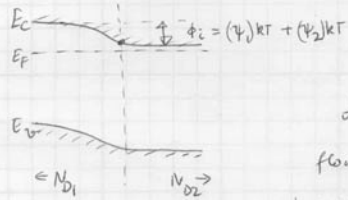
**PROBLEM 4 SOLUTION**

ASSIGNMENT 4 - EE 566 - Solid State Devices, Spring 2005. - (Solution)

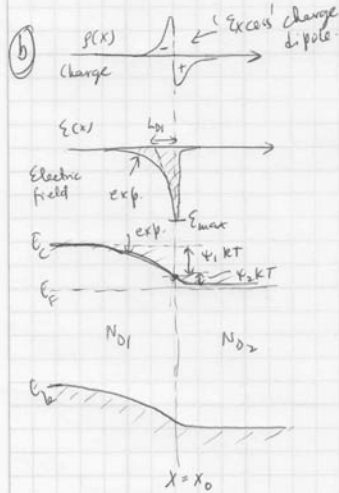
Problem 1 -

(a) Potential barrier is  $\phi_i = \frac{kT}{q} \ln\left(\frac{N_{D2}}{N_{D1}}\right) = 0.12 \text{ Volt}$

$\phi_i = 120 \text{ meV}$



for electron flow from region 1  $\rightarrow$  2, there is no barrier. However, for electron flow from 2  $\rightarrow$  1, there is a potential barrier  $\phi_i \approx 120 \text{ meV}$ . However, this barrier is very small, & any applied bias greater than 120 mV will make the current ohmic in both directions. So, there is rectification, but it is VERY WEAK.



(c) The exact solution to Poisson  $\nabla^2 \psi$  is

$$\frac{\partial^2 \psi(x)}{\partial x^2} = \frac{q N_D(x)}{\epsilon_s kT} [1 - e^{-q\psi(x)}]$$

which yields for Electric field (Don't class!)

$$|E(x)|^2 = \frac{2kT N_D}{\epsilon_s} [e^{q\psi(x)} + e^{-q\psi(x)} - 1]$$

Denoting total band bending in  $N_{D1}$  region as  $\psi_1 kT$  & in  $N_{D2}$  region as  $\psi_2 kT$ , and equating the field at  $x = x_0$ , the interface,

$$\left\{ \frac{2kT N_{D1}}{\epsilon_s} [-1 + e^{1\psi_1}] - 1 \right\} = \frac{2kT N_{D2}}{\epsilon_s} [1 + e^{-1\psi_2} - 1]$$

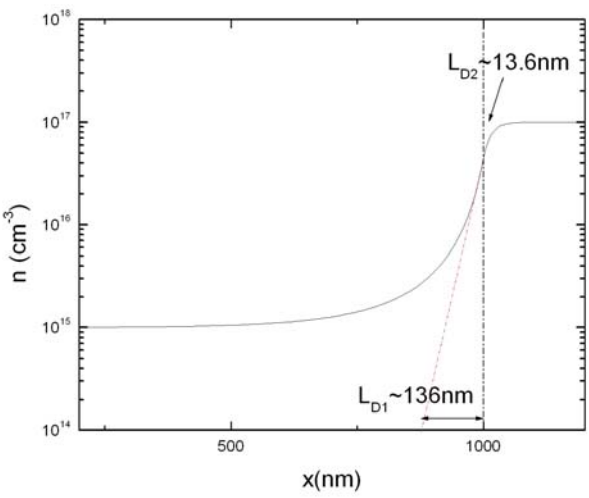
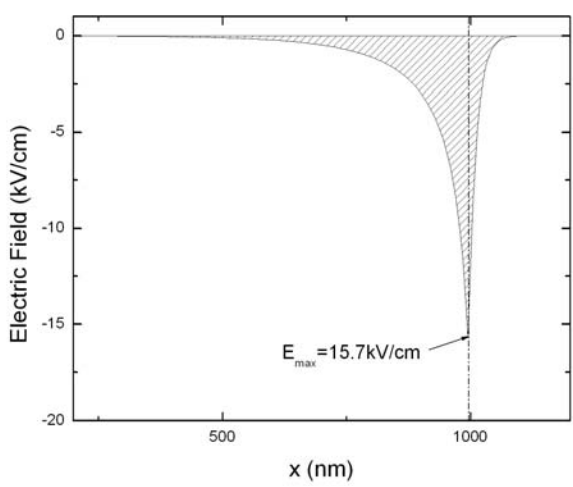
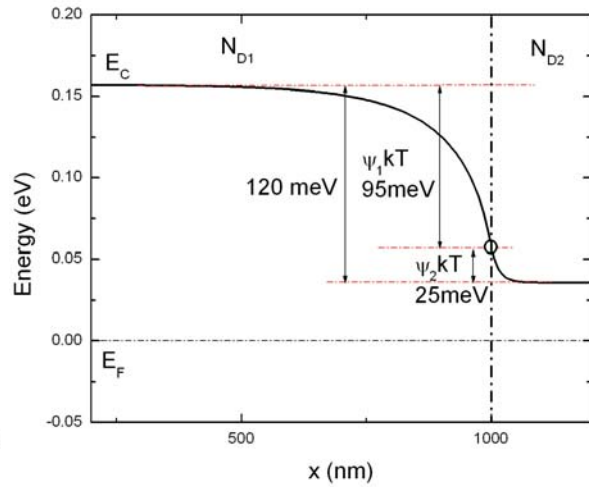
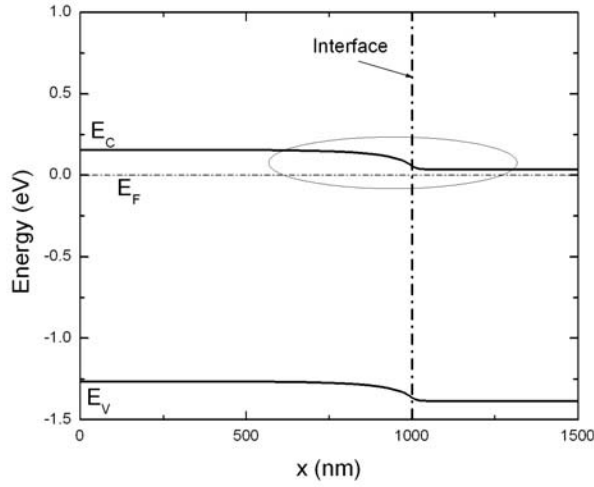
also,  $1\psi_1 + 1\psi_2 = \frac{\phi_i}{kT} = \frac{120 \text{ meV}}{26 \text{ meV}} @ 300 \text{ K}$ .

Solve simultaneously to get  $\boxed{1\psi_1 = 3.6517}$   
 $\boxed{1\psi_2 = 0.9636}$

$\Rightarrow$  45 meV drops in  $N_{D1}$ , & 25 meV drops in  $N_{D2}$ , &

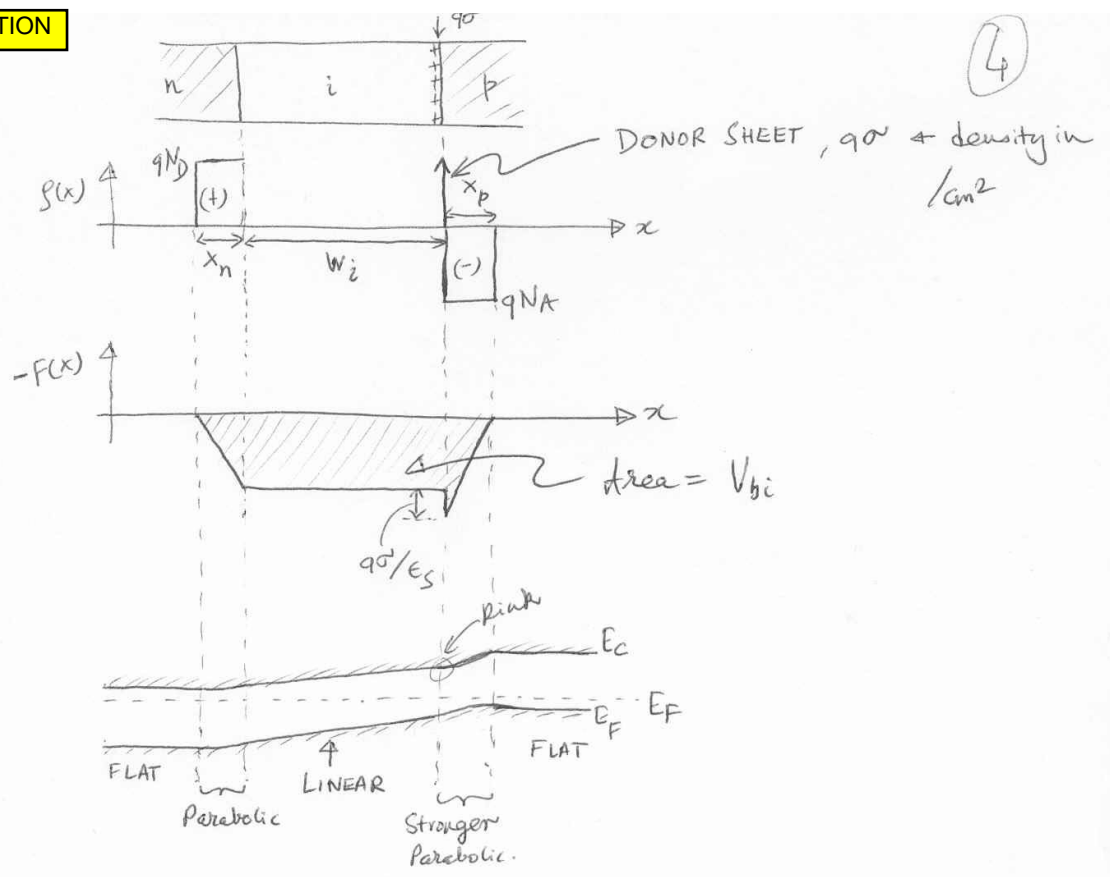
$$E_{max} (\text{exact}) = \left\{ \frac{2kT N_{D1}}{\epsilon_s} (-1 + e^{1\psi_1}) - 1 \right\}^{1/2} \approx \boxed{15.7 \text{ kV/cm}}$$

The simulated band diagrams using 1D Poisson are shown below. As can be seen, the band bending, the maximum electric field at the interface, and the Debye lengths match our exact calculations very well. This example illustrates that only for very simple structures can the Poisson equation be solved exactly. Most of the simulation for real devices is done numerically.



**PROBLEM 5 SOLUTION**

(c)



Charge Neutrality:  $qN_A x_p = qN_D x_n + q\sigma$

$\therefore x_p = x_n + \frac{\sigma}{N_A} \Rightarrow V_{bi} = \frac{qN_A x_p^2}{\epsilon_s} + \left( \frac{W_i qN_A}{\epsilon_s} - \frac{q\sigma}{\epsilon_s} \right) x_p$

$\left[ + \frac{1}{2} \frac{q\sigma^2}{\epsilon_s N_A} - \frac{W_i q\sigma}{\epsilon_s} \right]$

AREA UNDER  $F(x) - x$  plot.

$C_{p-i-n} = C_{p-n}$

(b)  $\frac{\epsilon_s}{x_n + x_p + W_i} = \frac{\epsilon_s}{W}$

$\uparrow$  with i-layer + sheet doping       $\uparrow$  No i-layer.

$\left( \frac{2\epsilon_s}{q} \cdot \frac{2}{N_A} \cdot V_{bi} \right)^{1/2}$

Solve (a) & (b) together to get -

$$W_i = (\sqrt{2} - 1) \frac{\sigma}{N_A} \Rightarrow W_i \approx (\sqrt{2} - 1) \times \frac{10 \times 5}{10^{17}} \text{ cm}$$

$W_i \approx 20.7 \text{ nm}$