

# EE566 Solid State Devices

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Dept of Electrical Engineering

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## Assignment 2 SOLUTIONS

EE 566 - ASSIGNMENT 2 Solutions

Fall 2005 Solid State Devices

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1. Assume the side of the 2D semiconductor is  $L$   
Then the density of states for 2D is

Ans →

$$g(E) = N(E)/E/L^2$$

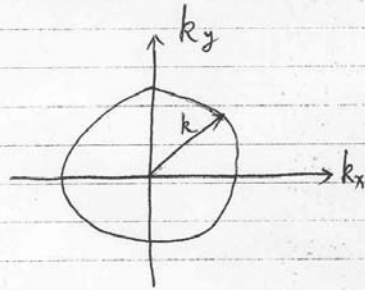
$$N(E) = 2 \cdot \frac{1}{4} \cdot \frac{\pi^2 k^2}{(\pi/L)^2}$$

$$= \frac{1}{2} \frac{k^2 L^2}{\pi}$$

$$E = \frac{p^2}{2m^*} = \frac{\hbar^2 k^2}{2m^*} \Rightarrow k^2 = \frac{2m^* E}{\hbar^2}$$

$$\therefore N(E) = \frac{m^* E L^2}{\pi \hbar^2}$$

$$\therefore g(E) = \frac{m^*}{\pi \hbar^2}$$



$$n_{2d} = \int_{E_c}^{\infty} g(E) f(E) dE$$

$$= \int_{E_c}^{\infty} \frac{m^*}{\pi \hbar^2} \cdot \frac{1}{1 + e^{(E-E_f)/kT}} dE$$

$$= \frac{m^*}{\pi \hbar^2} kT \int_{\frac{E_c - E_f}{kT}}^{\infty} \frac{1}{1 + e^{(E-E_f)/kT}} d\left(\frac{E-E_f}{kT}\right)$$

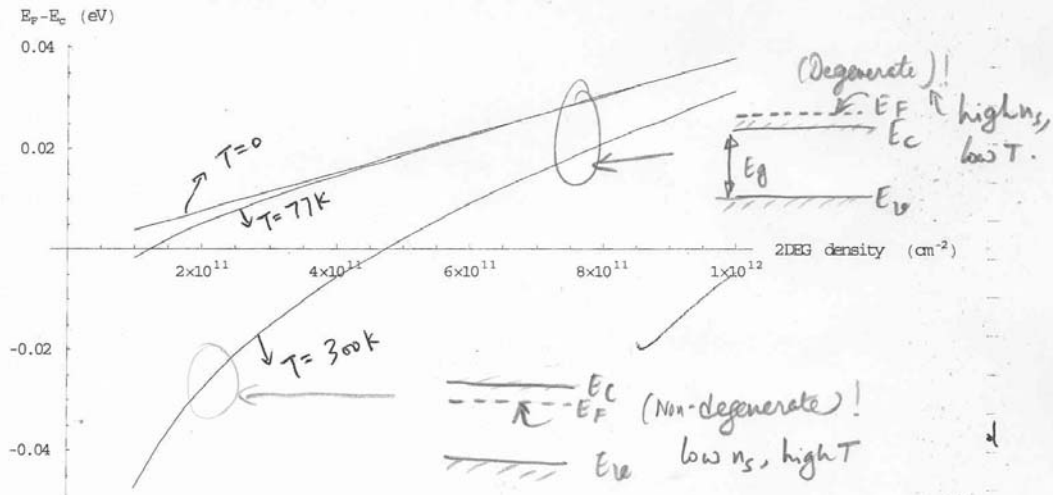
$$= \frac{m^*}{\pi \hbar^2} kT \cdot \ln\left(1 + e^{\frac{E_f - E_c}{kT}}\right)$$

$$\text{Let } N_c^{2d} = \frac{m^*}{\pi \hbar^2} kT, \quad \eta = \frac{E_f - E_c}{kT}$$

$$\text{Then } n_{2d} = N_c^{2d} (1 + e^{\eta})$$

$$\text{At } T=77\text{K}, 300\text{K}; E_F - E_c = kT \ln \left[ \exp \left( \frac{n_{2d} \pi \hbar^2}{m^* kT} \right) - 1 \right]$$

$$\text{At } T=0, E_F - E_c = \frac{n_{2d} \pi \hbar^2}{m^*}$$



$$n_{2d} = \frac{m^* kT}{\pi \hbar^2} \ln \left( 1 + e^{\frac{E_F - E_c}{kT}} \right)$$

From the above figure, it is shown that at low temperature the 2DEG carrier distribution is DEGENERATE.

$\therefore E_F > E_c$  and for small  $T$ ,  $\frac{E_F - E_c}{kT} \gg 1$

$$\text{then } n_{2d} = \frac{m^* kT}{\pi \hbar^2} \cdot \ln \left( e^{\frac{E_F - E_c}{kT}} \right) = \frac{m^*}{\pi \hbar^2} (E_F - E_c)$$

$\therefore$  at low temperature,  $n_{2d}$  is independent of temperature.

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So plot figure as follows. (fig. 1).

(c) 2DEG carrier distribution is degenerate

2) At room temperature, assume  $n = N_D(x)$

When equilibrium condition is satisfied,  $\bar{J} = 0$

$$\text{i.e. } \bar{J} = q n \mu \bar{E} + q D_n \frac{dn}{dx} = 0$$

$$\Rightarrow \bar{E}(x, T) = \frac{-D_n}{\mu n} \frac{dn}{dx} \\ = \frac{-D_n}{\mu N_D(x)} \cdot \frac{dN_D(x)}{dx}$$

Solution

Jing  $\rightarrow$

$$\text{From } \frac{D_n}{\mu} = \frac{kT}{q}$$

We get the relationship between  $\bar{E}$  and  $T$  is

$$\bar{E}(x, T) = - \frac{kT}{q N_D(x)} \frac{dN_D(x)}{dx}$$

b) As temperature increase, the magnitude of  $\bar{E}$  increases  
proportional to the temperature. The direction of  $\bar{E}$   
will not change. All we can say is that  $\bar{E} \uparrow$  as  $T \uparrow$ .

c) For constant doping  $\frac{dN_D(x)}{dx} = 0 \Rightarrow \bar{E}(x, T) = 0$   
 $\bar{E} = f(N_D, n)$  which change with temperature too.

$$d) \quad \bar{E} = - \frac{kT}{q N_D e^{-\frac{x}{\lambda}}} \cdot N_D e^{-\frac{x}{\lambda}} \left(1 - \frac{2x}{\lambda}\right) \\ = + \frac{2kTx}{q \lambda^2}$$

Please refer to fig. 2.

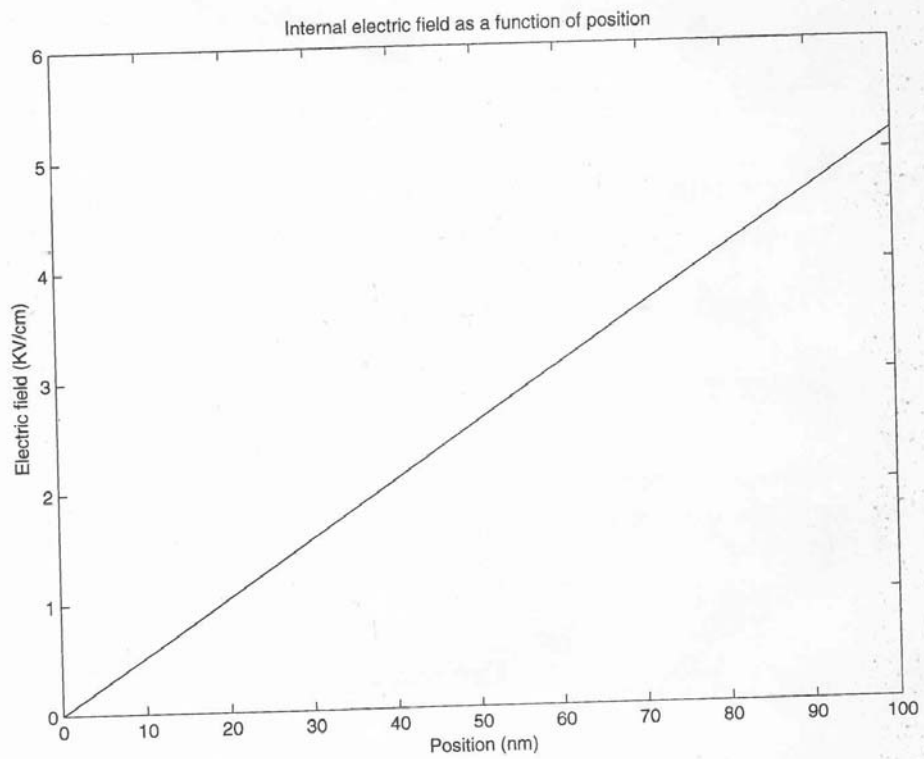
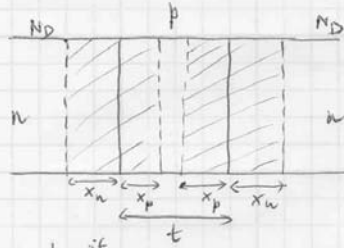
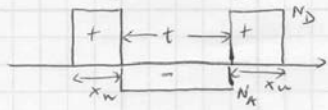


Fig 2.

Problem 3



if p-layer totally depleted



The simplest design will be to shoot for only ONE capacitor - this is achieved if the p-region is totally depleted!

∴ Requirement:  $2x_p \geq t$  — (1)

Charge neutrality →

$2x_n N_D = N_A t$

$x_n = \frac{1}{2} \frac{N_A \cdot t}{N_D}$  — (2)

22-141 50 SHEETS  
22-142 100 SHEETS  
22-144 200 SHEETS  
SAMPAD

Solution (He) →

If p-region doping is \$N\_A\$, the depletion width \$W\$ of a 'single' p-n junction is given by -

$V_{bc} = \frac{kT}{q} \ln \left( \frac{N_A N_D}{n_i^2} \right) = \frac{q}{2\epsilon_s} \left( \frac{N_A N_D}{N_A + N_D} \right) W^2$  — (3) = 0.84 Volt

From (1) + (2),  $(x_n + x_p)^2 = W^2 = \left( 1 + \frac{N_A}{N_D} \right)^2 \frac{t^2}{4}$  — (4)

Combining (3) + (4), we get, by eliminating \$W\$, of using  $d = 2x_n + t = t \left( 1 + \frac{N_A}{N_D} \right)$ ,

$\frac{1}{N_A} + \frac{1}{N_D} > \frac{q^2 d^2}{8\epsilon_s RT \ln \left( \frac{N_A N_D}{n_i^2} \right)}$  — (5)

∴  $N_A < N_A^{cr}$ , where  $N_A^{cr}$  is the solution of (5).

For  $N_D = 10^{17}/cm^3$  +  $d = 300 nm$ , I get

$N_A \leq 9.3 \times 10^{16}/cm^3$  if 't' scales accordingly.

For exactly  $\epsilon = \frac{\epsilon_s}{300nm}$  @  $V_a = 0$ ,  $N_A = 9.3 \times 10^{16}/cm^3$   $t = 2x_p^{cr} = 2 \times \frac{W \leq 150nm}{\left( 1 + \frac{N_A}{N_D} \right)} = 2 \times (78nm)$

⇒  $t = 156 nm$  will work!

Check:  $x_p = \frac{W}{1 + \frac{N_A}{N_D}} = 78nm$   $x_n = 72nm$  ⇒  $W = 2(x_n + x_p) = 2x_n + t = 300nm$