

# Acoustic Waves in a Circular Duct

Consider a circular duct of radius  $a$ . We take a cylindrical coordinate system  $\{x, r, \theta\}$ , where the  $x$  axis is along the duct axis. The acoustic pressure is governed by the wave equation

$$\frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} - \nabla^2 p = 0. \quad (1)$$

The pressure must satisfy an initial condition at  $x = x_0$  and a wall boundary condition at  $r = a$ . We use the method of separation of variables and assume

$$p(x, r, \theta, t) = X(x)R(r)\Theta(\theta)T(t). \quad (2)$$

Substituting (2) into (1) and dividing by  $X(x)R(r)\Theta(\theta)T(t)$ , gives

$$\frac{X''}{X} + \frac{R'' + R'/r}{R} + \frac{\Theta''}{r^2\Theta} - \frac{1}{c^2} \frac{T''}{T} = 0. \quad (3)$$

If we take

$$\frac{\Theta''}{\Theta} = -m^2, \quad (4)$$

$$\frac{X''}{X} = -k^2, \quad (5)$$

$$\frac{T''}{T} = -\omega^2, \quad (6)$$

where  $m$  is an integer. This implies a solution of the form

$$p_{mk\omega} = R_m(r)e^{i(kx+m\theta-\omega t)}. \quad (7)$$

The function  $R_m$  satisfies the equation

$$r^2 R_m'' + r R_m' + (\mu^2 r^2 - m^2) R_m = 0, \quad (8)$$

where we have introduced the eigenvalue  $\mu^2 = \omega^2/c^2 - k^2$ . For a rigid duct, this equation must satisfy an impermeability condition

$$\left(\frac{dR_m}{dr}\right)_{r=a} = 0. \quad (9)$$

Introducing the non-dimensional variable  $\tilde{r} = \mu r$ , equation (9) becomes

$$\tilde{r}^2 \frac{d^2 R_m}{d\tilde{r}^2} + \tilde{r} \frac{dR_m}{d\tilde{r}} + (\tilde{r}^2 - m^2) R_m = 0. \quad (10)$$

We recognize the Bessel equation and since the pressure is finite along the axis  $R_m = J_m(\tilde{r})$ . The wall condition (9) implies

$$J_m'(\mu a) = 0. \quad (11)$$

The boundary-value problem (10, 11) is a Sturm-Liouville problem whose solutions form a complete set. The derivative of the Bessel function has an infinite number of zeros which we denote as  $\{\alpha_{mn}\}$ ,

$$J'_m(\alpha_{mn}) = 0, \quad m = 0, 1, \dots \quad (12)$$

Hence, the eigenvalues are

$$\mu_{mn} = \frac{\alpha_{mn}}{a}. \quad (13)$$

This defines the axial wave number as

$$k_{mn} = \sqrt{\left(\frac{\omega}{c}\right)^2 - \mu_{mn}^2}. \quad (14)$$

The eigenfunction

$$p_{mn} = J_m\left(\frac{\alpha_{mn}r}{a}\right)e^{i(k_{mn}x+m\theta-\omega t)} \quad (15)$$

is called the  $\{mn\}$  mode. For every frequency  $\omega$ , the solution is then

$$p_\omega = \sum_{m=-\infty}^{m=+\infty} \sum_{n=0}^{n=+\infty} c_{mn} p_{mn}. \quad (16)$$

The expression for the coefficients  $c_{mn}$  is determined using the initial condition

$$p_\omega(0, r, \theta, t) = f_\omega(r, \theta)e^{-i\omega t} \quad (17)$$

and the orthogonality of the Bessel functions,

$$c_{mn} = \frac{1}{\pi a^2} \frac{\alpha_{mn}^2}{(\alpha_{mn}^2 - m^2) J_m^2(\alpha_{mn})} \int_0^{2\pi} \int_0^a f_\omega(r, \theta) J_m\left(\alpha_{mn} \frac{r}{a}\right) e^{-im\theta} r dr d\theta, \quad (18)$$

where we have used (see Hildebrand, p. 229)

$$\int_0^a r J_m^2\left(\alpha_{mn} \frac{r}{a}\right) dr = \frac{a^2(\alpha_{mn}^2 - m^2)}{2\alpha_{mn}^2} J_m^2(\alpha_{mn}). \quad (19)$$

Note the condition for propagation of an acoustic mode is that the wave number  $k_{mn}$  must be real. Otherwise the wave will decay exponentially and is known as an evanescent wave. Therefore an  $\{mn\}$  mode propagates if

$$\frac{\omega a}{c} > \alpha_{mn}. \quad (20)$$

At low frequencies, only the fundamental mode

$$p_{00} = e^{i[(\omega/c)x - \omega t]} \quad (21)$$

propagates. As  $\omega$  increases an additional mode propagates. The frequency at which a mode  $\{mn\}$  begins to propagate is known as the cutoff frequency of the mode. As the frequency increases (decreases) and is equal to the cutoff frequency of a mode  $\{mn\}$ , the mode  $\{mn\}$  is said to cut on (cut off).

As an example, consider a duct of radius  $a = 0.5m$ ,  $c = 340m/s$ , and the sound frequency is  $3000rpm$ .  $\omega/c = 0.462$ . From the tables of zeros of Bessel functions, the lowest zero is  $\alpha_{11} = 1.8412$ . hence only the fundamental mode will propagate.

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## Bessel Function Zeros

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The first  $k$  roots  $x_1, \dots, x_k$  of the Bessel function  $J_\alpha(x)$  are given in the following table. They can be found in *Mathematica* using the command `BesselJZeros[n, k]` in the *Mathematica* add-on package `NumericalMath`BesselZeros`` (which can be loaded with the command `<<NumericalMath``).

zero	$J_0(x)$	$J_1(x)$	$J_2(x)$	$J_3(x)$	$J_4(x)$	$J_5(x)$
1	2.4048	3.8317	5.1356	6.3802	7.5883	8.7715
2	5.5201	7.0156	8.4172	9.7610	11.0647	12.3386
3	8.6537	10.1735	11.6198	13.0152	14.3725	15.7002
4	11.7915	13.3237	14.7960	16.2235	17.6160	18.9801
5	14.9309	16.4706	17.9598	19.4094	20.8269	22.2178

The first  $k$  roots  $x_1, \dots, x_k$  of the derivative of the Bessel function  $J'_\alpha(x)$  can be found in *Mathematica* using the command `BesselJPrimeZeros[n, k]` in the *Mathematica* add-on package `NumericalMath`BesselZeros`` (which can be loaded with the command `<<NumericalMath``). The first few such roots are given in the following table.

zero	$J'_0(x)$	$J'_1(x)$	$J'_2(x)$	$J'_3(x)$	$J'_4(x)$	$J'_5(x)$
1	3.8317	1.8412	3.0542	4.2012	5.3175	6.4156
2	7.0156	5.3314	6.7061	8.0152	9.2824	10.5199
3	10.1735	8.5363	9.9695	11.3459	12.6819	13.9872
4	13.3237	11.7060	13.1704	14.5858	15.9641	17.3128
5	16.4706	14.8636	16.3475	17.7887	19.1960	20.5755

**SEE ALSO:** [Bessel Function](#), [Bessel Function of the First Kind](#). [[Pages Linking Here](#)]

## CITE THIS AS:

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